

Ch 5 Lattice Patterns



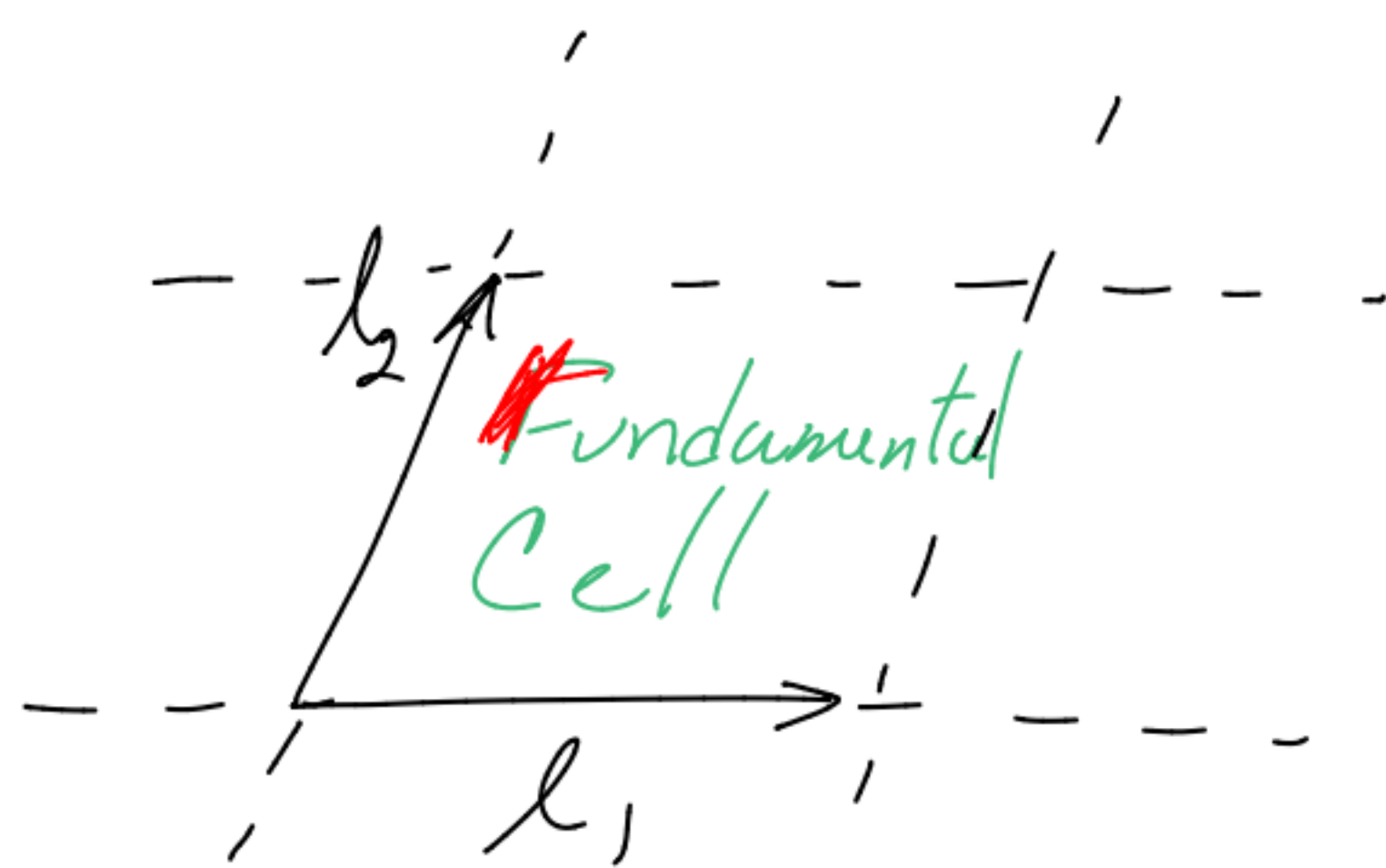
Heat

Let $\vec{l}_1$ & $\vec{l}_2$ be two lin.
ndep. vectors in $\mathbb{R}^2$

Suppose: $U(x, t): \mathbb{R}^2 \rightarrow \mathbb{R}$

U is doubly periodic w.r.t.

$\vec{l}_1$ and $\vec{l}_2$ if $U(x + \vec{l}_1, t) = U(x, t)$
 $= U(x + \vec{l}_2, t)$



Let $\vec{l} = n_1 \vec{l}_1 + n_2 \vec{l}_2, n_1, n_2 \in \mathbb{Z}$

$U(x + \vec{l}, t) \stackrel{\text{want}}{=} U(x, t)$

Def. Let $\mathcal{L} = \{ \vec{l} = n_1 \vec{l}_1 + n_2 \vec{l}_2 : n_1, n_2 \in \mathbb{Z} \}$
planar lattice (infinite discrete group)
of $\mathbb{R}^2$

Def $C_L^\infty = \left\{ \text{All smooth functions} \right.$
 $\left. \text{that are doubly per.} \right.$
 $\text{w.r.t. } L$

$$C_L^\infty = \left\{ u \in C_L^\infty : u(x+l, t) = u(x, t), \right. \\ \left. \forall l \in L \right\}$$

Remarks

The lattice L is a subgroup of translations

C_L^∞ is the fixed-point subspace of the action of L on C_L^∞

There are five kinds of planar lattices: rhombic, square, hexagonal, rectangular and oblique

3. C_L^∞ is an inv. subspace

5. The lattices that are most relevant for bif. theory are those generated by two vectors of the same length: rhombic, square, hexagonal.

Def. $L^* = \{ n_1 \vec{k}_1 + n_2 \vec{k}_2 : n_1, n_2 \in \mathbb{Z} \}$
 $\vec{k}_1, \vec{k}_2 \in \mathbb{R}^2$ s.t.

$$\vec{k}_i \cdot \vec{l}_j = 2\pi \delta_{ij}, \quad i, j = 1, 2$$

OR

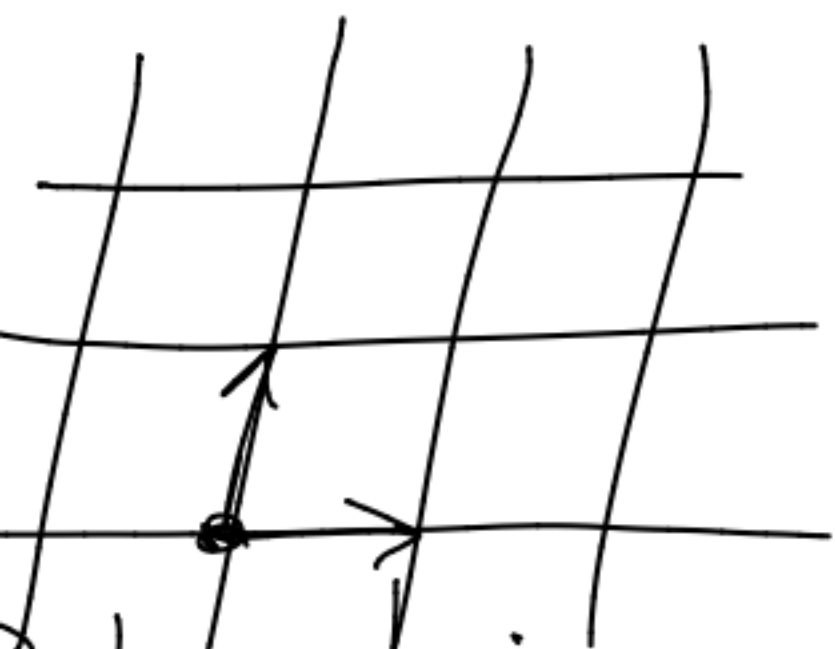
$$L^* = \{ \vec{k} \in \mathbb{R}^2 : \vec{k} \cdot \vec{l} \in \mathbb{Z} \{2\pi\}, \forall \vec{l} \in L \}$$

Consider $u(x,t)$ as a sum of Fourier modes that lie on the dual lattice:

$$u(x,t) = \sum_{\mathbf{k} \in \mathcal{L}^*} Z_{\mathbf{k}}(t) e^{i\mathbf{k} \cdot \mathbf{x}} + \text{c.c.}$$

$\underbrace{Z_{\mathbf{k}}(t)}_{\text{amplitude coeff.}} \underbrace{e^{i\mathbf{k} \cdot \mathbf{x}}}_{\text{spatial modes}}$

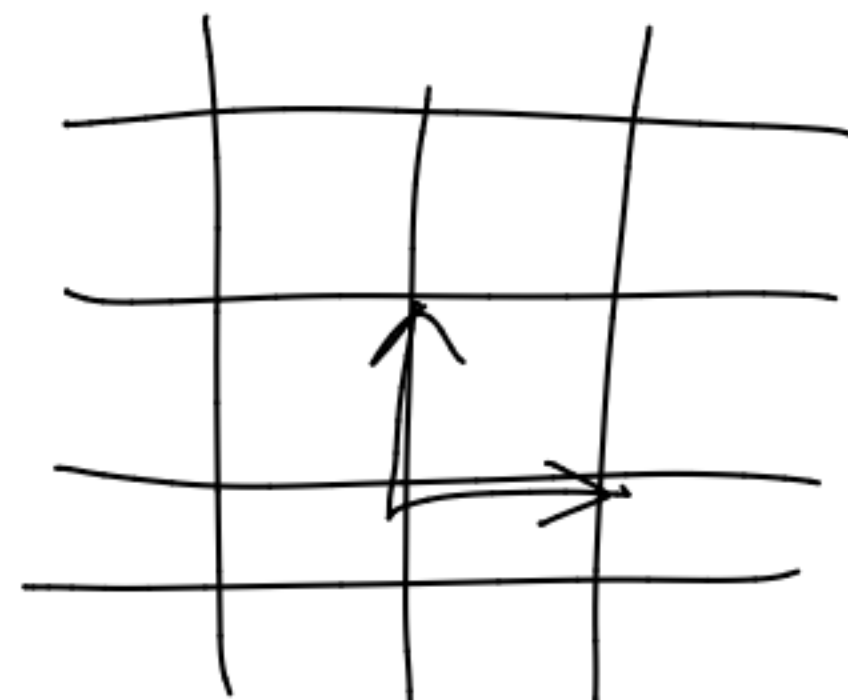
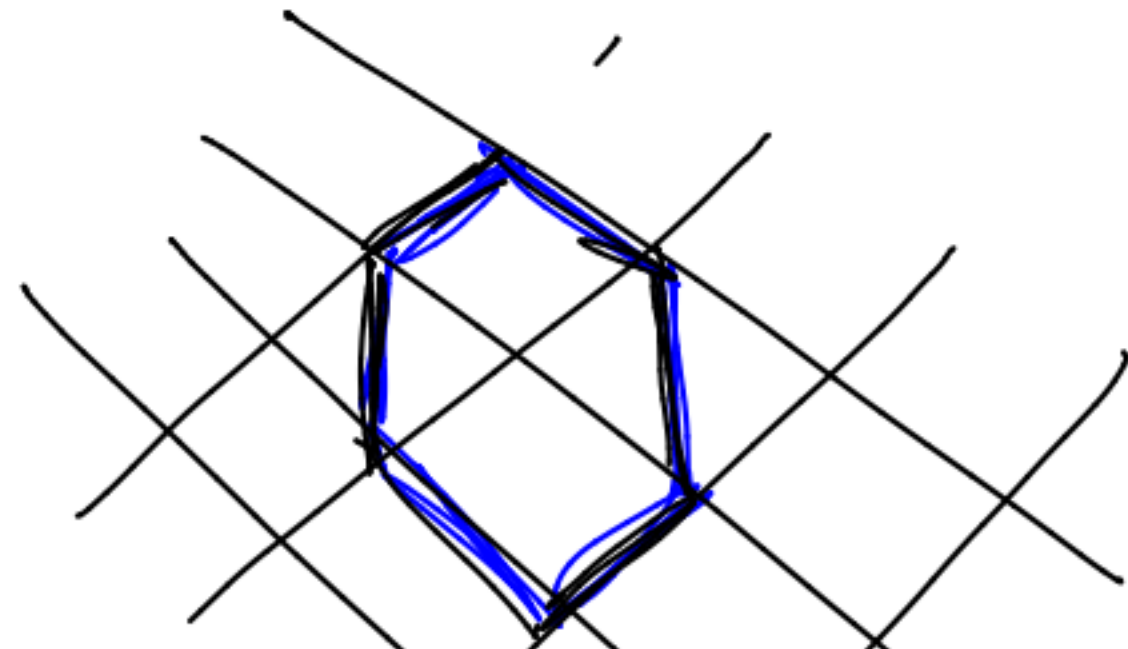
Common dual lattices:



Rhombic

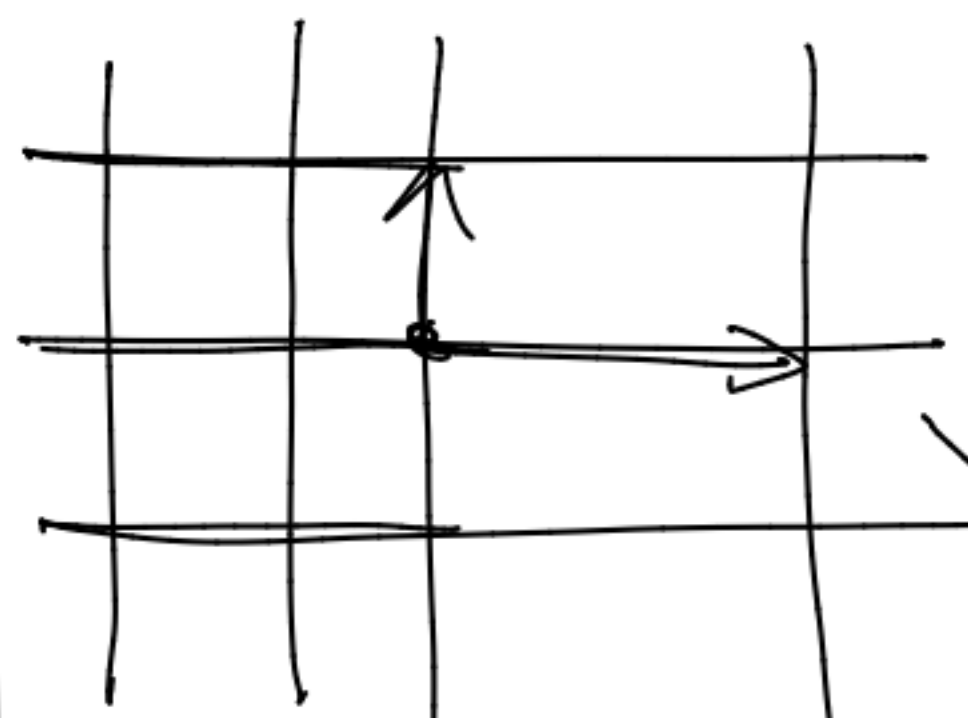
$$\vec{k}_1 = k^2 \cos\left(\frac{\pi}{3}\right), \vec{k}_2 = k^2 \cos\left(\frac{2\pi}{3}\right)$$

Hexagonal

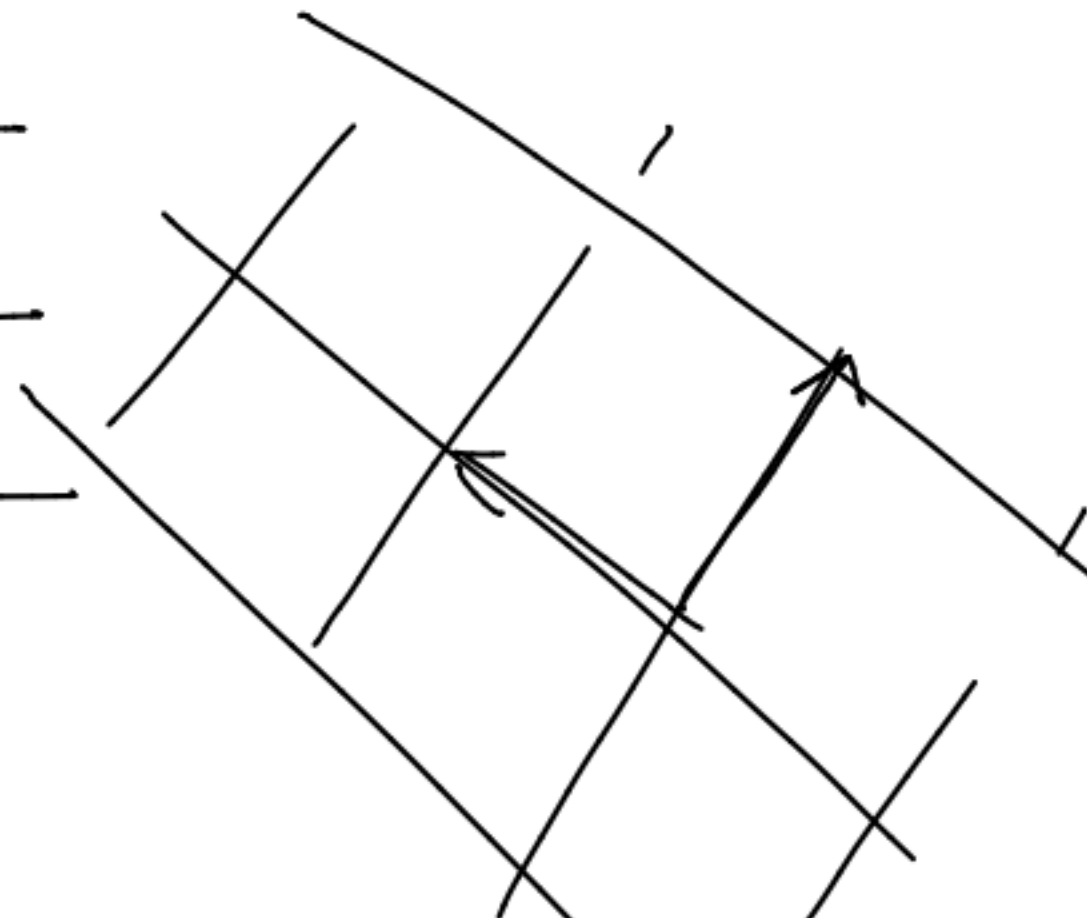


Squares

$$|\vec{k}_1| = |\vec{k}_2|$$



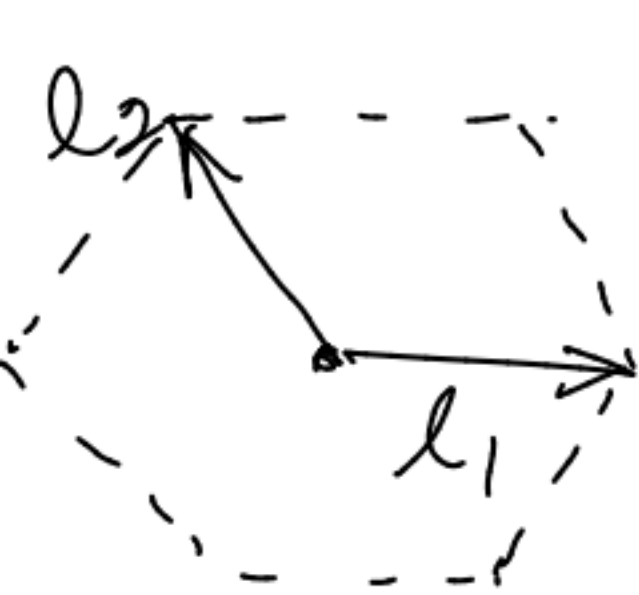
Rectangular



Centered Rectangular

g.) Hexagonal Lattice

$$l = \left\{ \vec{l}_1 = (c, 0), \vec{l}_2 = \left(-\frac{c}{2}, c\frac{\sqrt{3}}{2}\right) \right\}$$



Let $\vec{k}_1 = (a, 0)$

$$\vec{k}_1 \cdot \vec{l}_1 = 2\pi$$

$$\vec{k}_1 \cdot \vec{l}_2 = 0$$

$$\vec{k}_1 \cdot \vec{l}_1 = ac \stackrel{\text{set}}{=} 2\pi \Rightarrow a = \frac{2\pi}{c}$$

$$\vec{k}_2 \cdot \vec{l}_1 = 0 \Rightarrow \vec{k}_2 = \frac{2\pi}{c} \left(0, \frac{2}{\sqrt{3}}\right)$$

$$\vec{k}_2 \cdot \vec{l}_2 = 2\pi$$

$$\vec{k}_1 = \frac{2\pi}{c} \left(1, \frac{1}{\sqrt{3}}\right)$$

Bifurcations on a Lattice

$$\frac{2u(x,t)}{2t} = f(u(x,t), \mu)$$

$\mu \equiv \text{bif. par.}, x \in \mathbb{R}^2$
 $u(x,t) \in \mathbb{R}^1$

Symmetry Group: $\Gamma = E(2)$
 Euclidean group of rotations,
 reflections and translations
 on a plane.

$\Rightarrow f$ must be $E(2)$ equiv.

$$\gamma f(x, \mu) = f(\gamma u, \mu), \forall \gamma \in E(2)$$

$E(2)$

Translation $\equiv p \in \mathbb{R}^2$: $p \cdot U(x) = U(x-p)$

Rotation & reflection $\equiv \sigma \in O(2)$

$$\sigma \cdot U(x) = U(\sigma^{-1} x)$$

⑤ σ^{-1} is needed for the action of $E(2)$ on the space of functions to be a homomorphism.

Pf Need: $\Theta(\gamma\sigma) = \Theta(\gamma)\Theta(\sigma)$

$$\forall \gamma, \sigma \in E(2)$$

Let $\boxed{\Theta(\gamma) U(x) = U(\gamma^{-1} x), \forall \gamma \in E(2)}$

$$W(x) = \Theta(\sigma) U(x) = U(\sigma^{-1} x)$$

$$\Theta(\gamma)(\Theta(\sigma) U)(x) = \Theta(\gamma) W(x) = W(\gamma^{-1} x)$$

$$= U(\sigma^{-1} \gamma^{-1} x) = U((\gamma\sigma)^{-1} x)$$

$$= \Theta(\gamma\sigma) U(x), \forall \sigma, \gamma \in E(2)$$

⑥ inverses ensure group elements act in correct order. $\square$

at $l \in L$

$$p \cdot u(x+l) = u((x+l)-p) = u(x-p+l) \\ = u(x-p) = p \cdot u(x)$$

$\Rightarrow C_L^\infty$ is inv. under all translations

Symmetry $E(2)$

(PDE) $\frac{\partial}{\partial t} u(x,t) = f(u(x,t), \mu)$

{decaying/
growing
modes
Sol.

$$u(x,t) = \sum_{j=1}^n z_j(t) e^{i \vec{k}_j \cdot \vec{x}} + \text{c.c.}$$

$z_j(t) \in \mathbb{C}$, $|\vec{k}_j| = k$, = critical wave number

$\Downarrow$
Amp. Eq. $\frac{dz}{dt} = g(z, \mu)$

$$z = (z_1, \dots, z_r) \in \mathbb{C}^n$$

Stationary
Trivial Sol. $z=0 \Leftrightarrow g(0, \mu)=0$

Def $\mathcal{H}_L$ = Holohedry of L

$$= \{ \sigma \in O(2) : \sigma L = L \}$$

$\odot C_L^\infty$ is inv. under $\mathcal{H}_L$
 $\mathcal{H}_L$ is the group of symmetries that the Amp. Eq. retains from $E(2)$ in the original PDE.